

Take a look at the Compton equation

$$\lambda' - \lambda_0 = \frac{h}{m_e c} (1 - \cos \theta)$$

$\frac{h}{m_e c}$ has units of wavelength

$m_e c$ has units of momenta

recall

for a photon

$$E = hf$$

$$p = \frac{h}{\lambda}$$

It was postulated by de Broglie that matter too possesses these specific qualities

$$E = hf \quad \lambda = \frac{h}{p}$$

Matter has a wavelength? Crazy!
Consider the evidence at that time.

Bohr was able to calculate the energy splittings in an atom with a simple suggestion.

angular momentum is quantized

$$mvr = n(h/2\pi)$$

recall

$$E = KE + PE$$

2

$$= \frac{1}{2}mv^2 + \left(-\frac{kq^2}{r}\right)$$

$$\frac{mv^2}{r} = \frac{kq^2}{r^2}$$

$$mv^2 = \frac{kq^2}{r}$$

$$E = -\frac{1}{2} \frac{kq^2}{r}$$

$$v^2 = \left(\frac{n}{mr} \left(\frac{h}{2\pi}\right)\right)^2 = \frac{kq^2}{mr}$$

$$\Rightarrow r = \frac{n^2 \left(\frac{h}{2\pi}\right)^2}{mkq^2} = n^2 (0.0529 \text{ nm})$$

$$E = -13.6 \text{ eV}/n^2 \quad E_{\text{photon}} = -13.6 \text{ eV} \left(\frac{1}{n_1^2} - \frac{1}{n_2^2}\right)$$

look again

$$mvr = n \frac{h}{2\pi}$$

$$2\pi r = n \frac{h}{mv} = n \frac{h}{p} = n\lambda$$

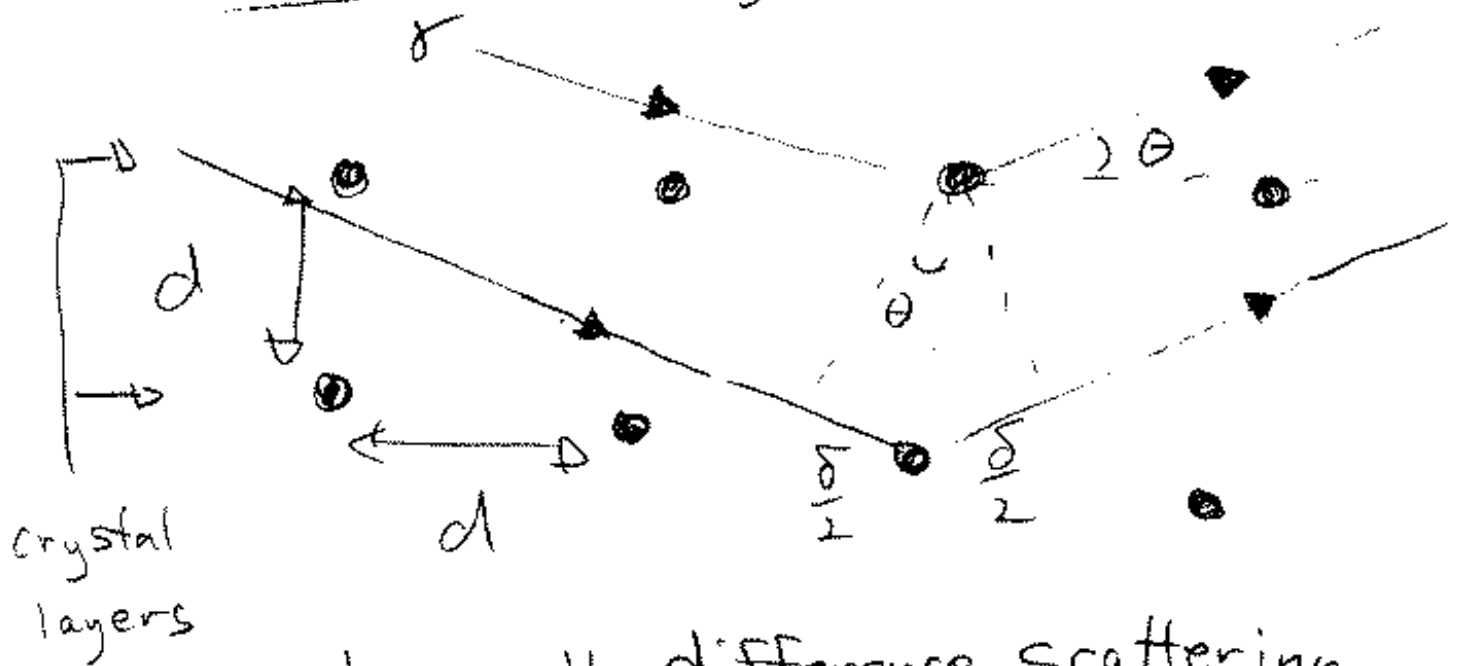
Bohrs condition can be interpreted as how many λ 's can safely sit in an orbit (standing wave condition)

lucky guess!

The confirmation arrived by accident 3

In the Compton effect, the λ of the emitted light was measured by looking at the diffraction angle of max intensity

Consider a crystal Lattice



get a path difference scattering
off of different layers

$$\delta = 2d \sin \theta$$

maxima when $2d \sin \theta = m\lambda$

Bragg
diffraction

2 experimenters were studying electron scattering from nickel when their sample got contaminated.

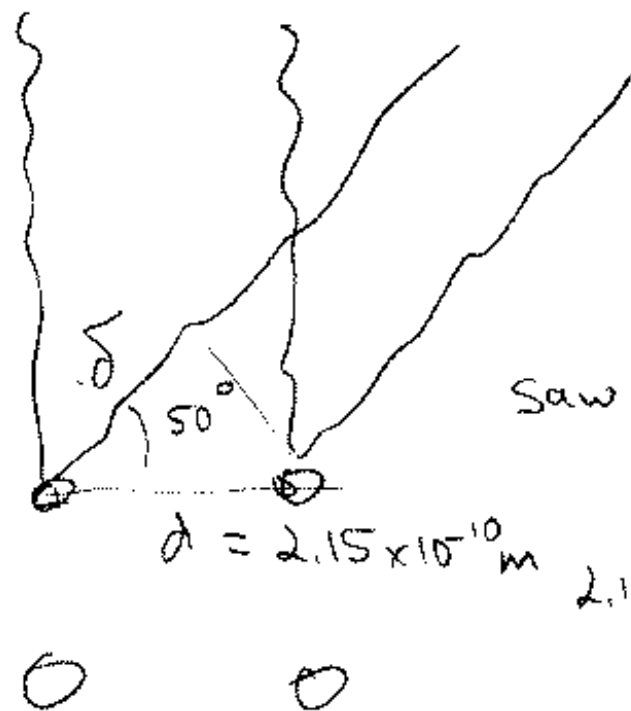
They cleaned it up by heating it up. They also made the nickel a crystal, so when they turned their electron beam back on:

Electrons accelerated through 54V have a KE of 54eV according to de Broglie

$$p = \sqrt{2m_e KE} \quad \lambda = \frac{hc}{\sqrt{2m_e c^2 KE}}$$

$$= \frac{1240 \text{ eV nm}}{\sqrt{2(511000 \text{ eV}) 54 \text{ eV}}}$$

$$= 0.167 \text{ nm}$$



Saw a peak @ 50°

$$d \sin \theta = \lambda \quad (\text{1st max})$$

$$2.15 \times 10^{-10} \text{ m} \sin 50^\circ = 1.65 \text{ nm}$$

pretty close

People were convinced something interesting was going on, but still had doubts.

ex get diffraction with electrons, but, how about one electron at a time through 2 slits?

$\Rightarrow$ Turned out you got the same pattern!

(cool discussion in Time line)
actually

You can think about the electron ^{or photon} existing as probabilities

ex Photon & 2 slits

a photon's probability function consists of the probability it goes through one slit & probability it goes through another

$$\psi(x, t) = \psi_1(x, t) + \psi_2(x, t)$$

↑
wave function

"wave" or
Probability
functions

$$= A_1 \sin(kx - \omega t + \theta_1) + A_2 \sin(kx - \omega t + \theta_2)$$

$$A_1 = A_2$$

$$\sum |\psi|^2 = 1 \quad \text{all probabilities add to one}$$

now, what we see at the far screen

$$\psi = A \sin(\theta_1) + A \sin(\theta_1 + \phi)$$

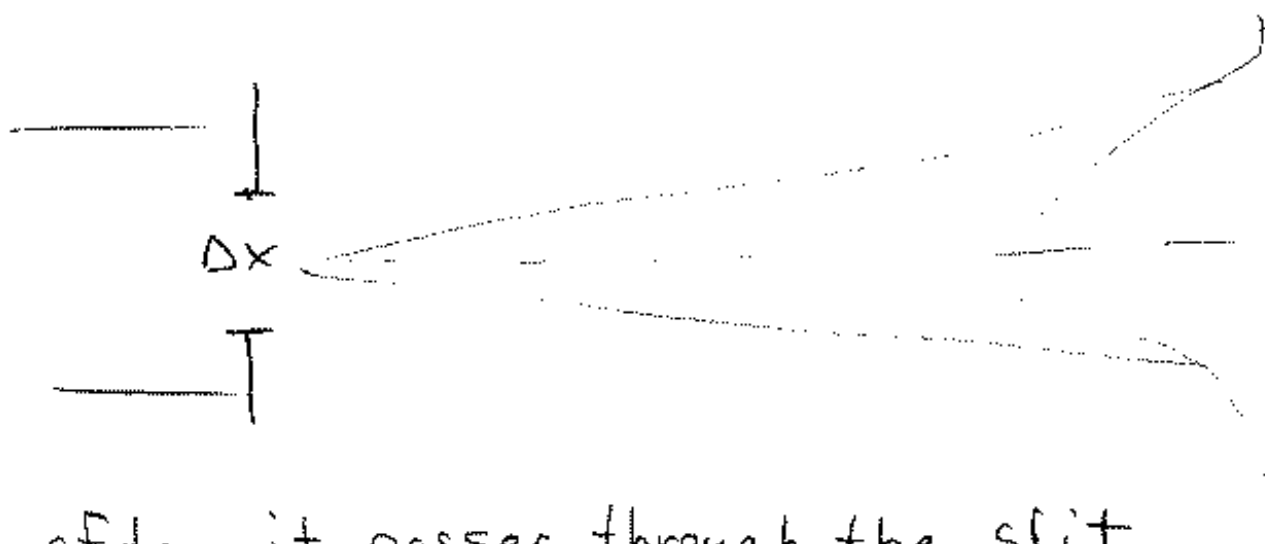
$$|\psi|^2 = A^2 \left(2 \cos \frac{\phi}{2} \sin \left(\theta_1 + \frac{\phi}{2} \right) \right)^2$$

as before

this has some bizarre
consequences.

I mean, fine, we get a 2 slit interference, 7
but if we cover one of the slits, we'll
surely see a single slit pattern. But
lets interpret it differently

as the wave passes through the slit,
we are forcing the wave to exist only
in the area of the slit



after it passes through the slit
we expect a minima at $\Delta x \sin \theta = \lambda$

might say that if we let in single
"particle" we can measure their location only
to within the location of the 1st peak
or we give the particle some uncertainty
in momentum $\uparrow$ when we confine it.

means it is difficult to know position & location arbitrarily well. 8

now, if we look at many "events" or say electrons striking a screen after a slit, we can do a pretty good job of finding the center, but, event by event we can't predict where the electron will fall within a given precision.

Consider a photon

$$\Delta x \Delta p_x \geq \frac{h}{4\pi} \quad E = pc$$

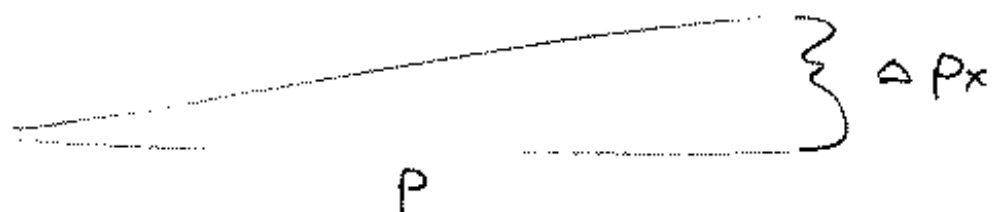
$$\frac{\Delta x}{c} \Delta p_x c \geq \frac{h}{4\pi}$$

$$\Delta t \Delta E \geq \frac{h}{4\pi}$$

now, think about a logic pulse, a fast one



when you work with these you notice you get a lot of noise on sensitive equipment



$$\frac{\Delta p_x}{p} \sim \theta$$

$$\Delta x \left(\frac{\Delta p_x}{p} \right) \sim \lambda$$

$$\Delta x \frac{\Delta p_x}{p} \sim \frac{h}{p}$$

$$\Delta x \Delta p_x \sim h$$

as we squeeze Δx (make it smaller)
 Δp_x compensates

~~consider a photon $E = pc$~~

$$\frac{\Delta x}{c} \Delta pc \sim \frac{h}{c}$$

Actually this is a rough form
 of the uncertainty principle

$$\Delta x \Delta p_x \geq \frac{h}{4\pi} \quad \text{more precisely}$$

It's because there are a lot of different frequencies going into making that sharp pulse

$\Rightarrow$ you know where an edge is extremely well, so you induce a spread of Δp if you like

or Δt is known very well

$\Delta E = h \Delta f$ is big

$$\Delta t (h \Delta f) \geq \frac{h}{4\pi}$$

$$\Delta t \Delta \omega \geq \frac{1}{2}$$

We can see this by looking at a small portion of the x axis when we add together some cosines

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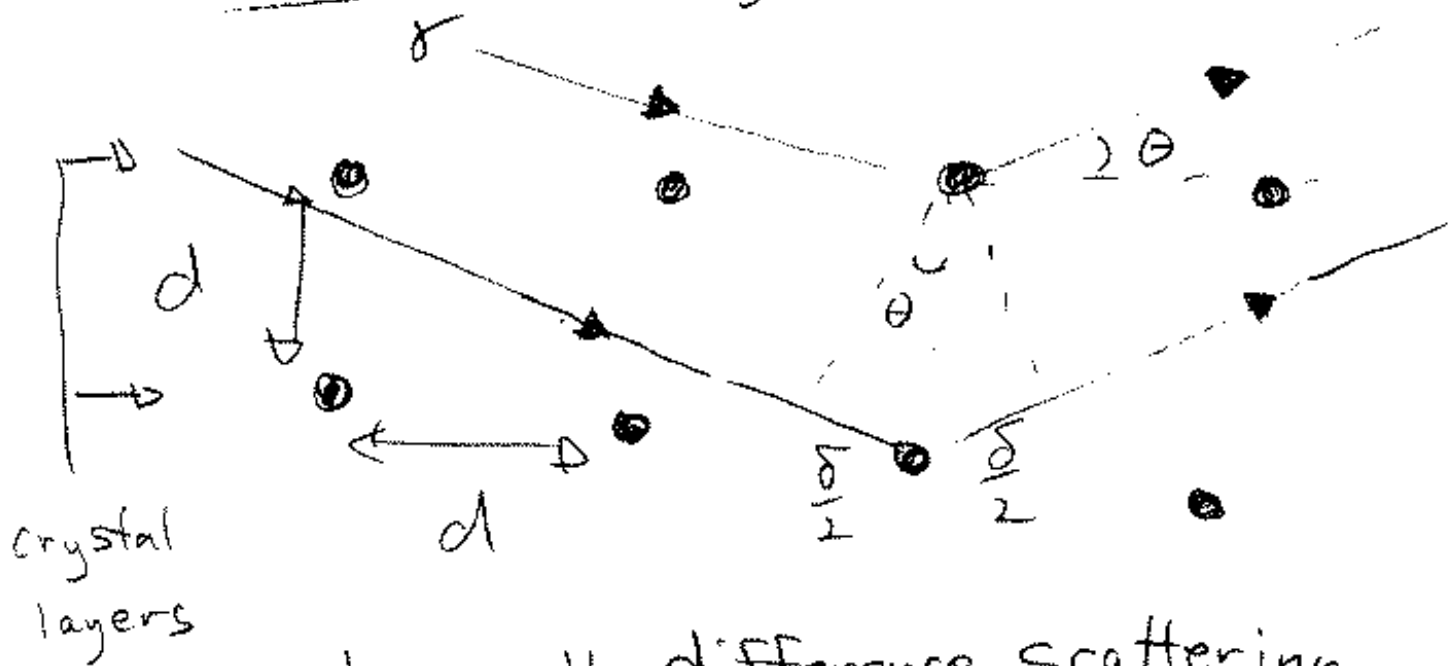
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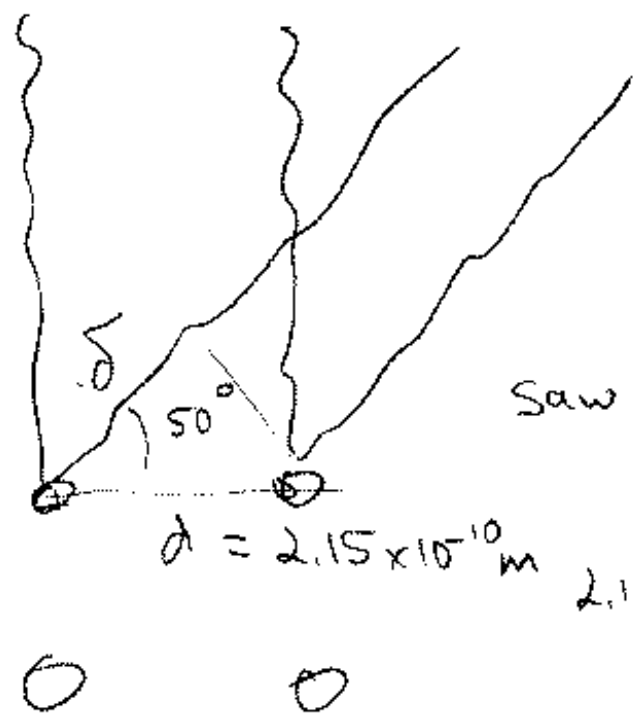
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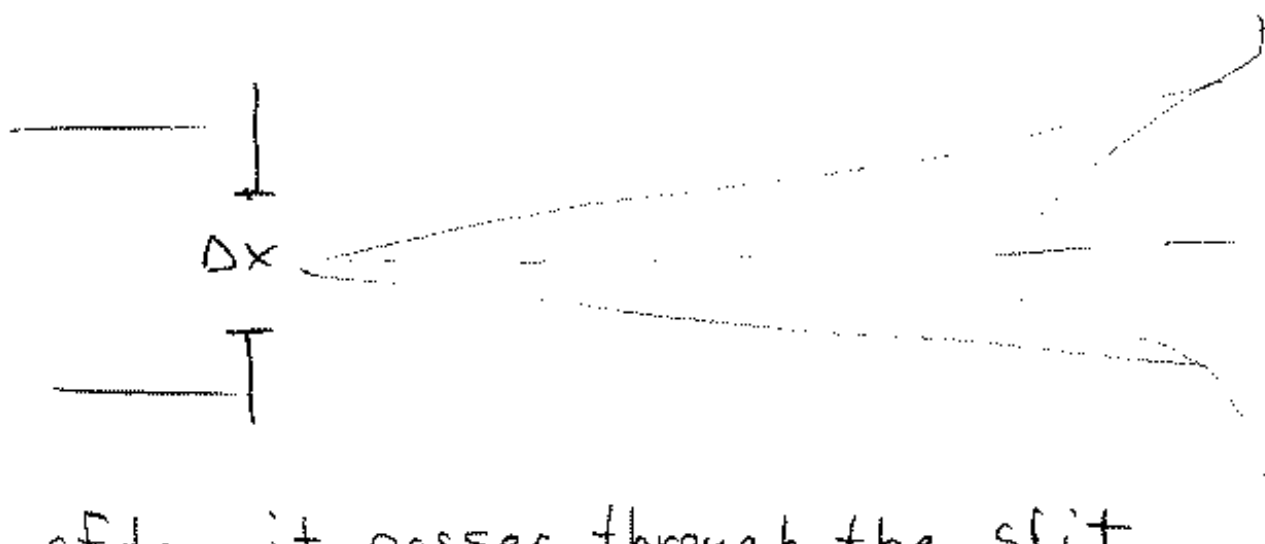
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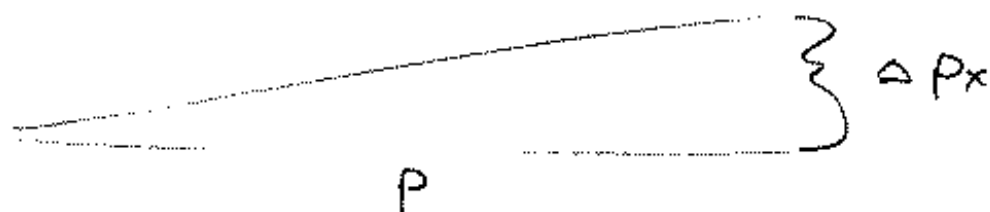
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