

$V=?$

know the electric field every where  
how do we get potential

lets make  $V_I = r$  inside

$V_F = \infty$

and well take  $V_F = 0$

to get to  $\infty$  we'll take a charge  
a charge from  $a \rightarrow b$ , then from  $b \rightarrow \infty$

$$V_F - V_I = (V_\infty - V_b) + (V_b - V_a) + (V_a - V_r)$$

$$= - \int_b^\infty \vec{E} \cdot d\vec{s} - \int_a^b \vec{E} \cdot d\vec{s} - \int_r^a \vec{E} \cdot d\vec{s}$$

$$= - \int_b^\infty \frac{kq_b}{r^2} dr - \int_a^b 0 \cdot ds - \int_r^a \frac{kq_b}{r^2} dr$$

$$= - \left[ -\frac{kq_b}{r} \Big|_b^\infty \right] - 0 - \left[ -\frac{kq_b}{r} \Big|_r^a \right]$$

$$-V_I = \left( -\frac{kq_b}{b} + \frac{kq_b}{a} - \frac{kq_b}{r} \right)$$

$$V_{\text{inside}} = \left( \frac{kq_b}{r} - \frac{kq_b}{a} + \frac{kq_b}{b} \right)$$

Lets continue our discussion of Gauss's Law LECTURE 3  
 (the relationship between surface shapes, Electric fields & charge)  
 Sept 7

$$\oint \vec{E} \cdot d\vec{A} = q_{\text{enclosed}} / \epsilon_0$$

Our goal is to connect it to something real. How do we make E?

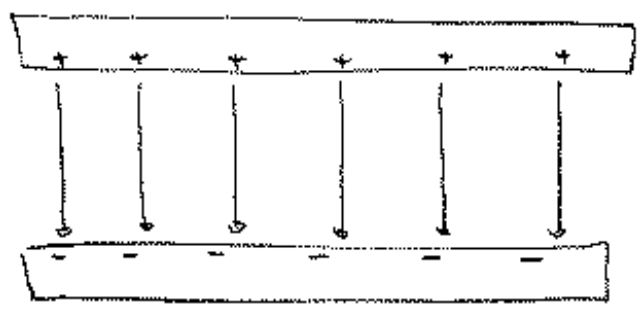
Consider 2 conducting sheets separated by a distance, d

point charge  $\vec{E} = \frac{q}{4\pi r^2}$

line of charge  $\vec{E} = \frac{\lambda}{2\pi r \epsilon_0}$

sheet of charge  $\vec{E} = \frac{\sigma}{2\epsilon_0}$

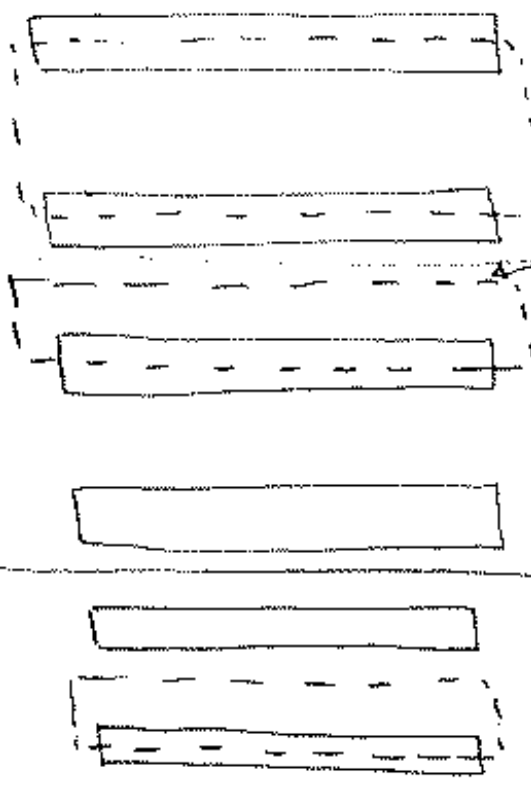
In between the sheets the field looks like



On the top sheet I have a charge +q  
 on the bottom sheet I have a charge -q

What can we glean from Gauss's Law (I see many ways)  
 - Choose a shape that mimics the E field (box)

$q_{\text{enc}} = 0$   
 $E = 0$  inside a conductor



$\vec{E} \cdot \hat{n} = 0$  at sides  
 $0 = 0$

Same story

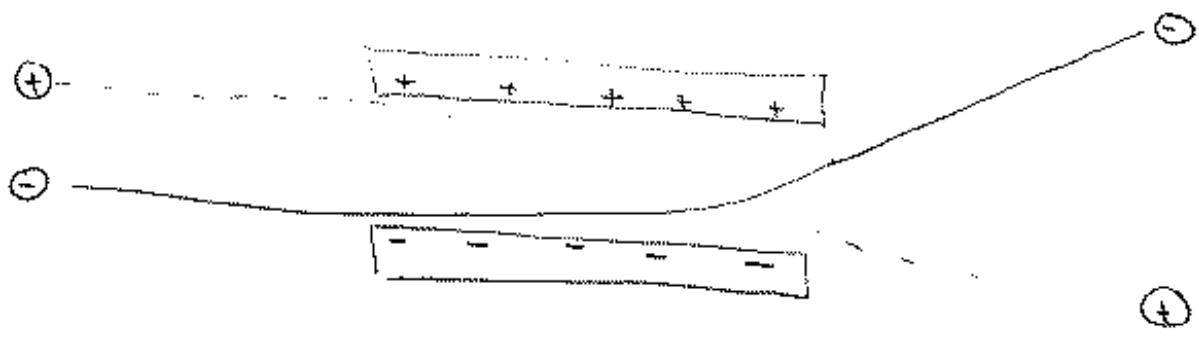
$E$  must be 0 here  
 Since  $q_{\text{enc}} = 0$   
 $E = 0$  in cond  
 $\vec{E} \cdot \hat{n} = 0$  outside

$E A = \frac{q_{\text{enc}}}{\epsilon_0}$   
 $E = \frac{q_{\text{enc}}}{A \epsilon_0}$   
 so what?

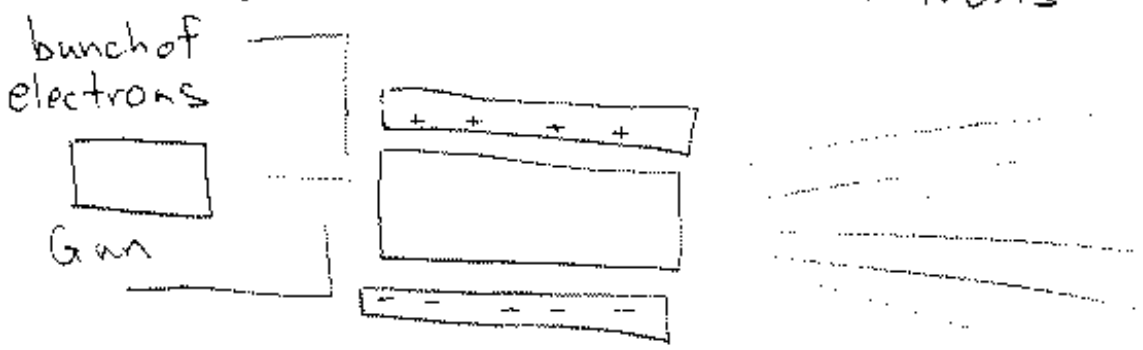
$q_{\text{enc}} = -q$   
 $\vec{E} \cdot \hat{n} = -E$   
 $E_x^2 = \frac{q_{\text{enc}}}{\epsilon_0}$   
 $A = x^2$

Check this out

What do you suppose happens if I shoot an electron between the plates? It bends up!



Now, imagine a whole bunch of electrons



and sheets on front back top & bottom that you could control the amount of charge on

Could be useful?

and a way to record the spot where the electron hit

Usually you see these things in objects like this

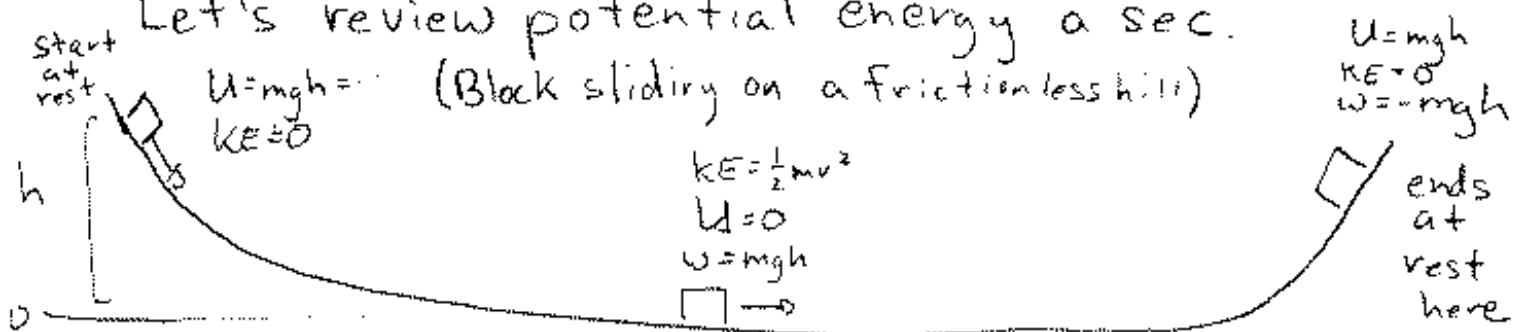


a TV (more likely an oscilloscope)

So that Electric field can be pretty handy some times

You notice that, in between the plates, (3)  
 work is being done on the charge. The  
 charge is moving from a place with high  
 potential energy to a place with lower  
 potential energy because the electric field is  
 doing work on it.

Let's review potential energy a sec.



$$\begin{aligned}\Delta U &= U_f - U_i \\ &= 0 - mgh \\ W &= mgh \\ \Delta U &= -W\end{aligned}$$

$$\begin{aligned}\Delta U &= U_f - U_i \\ &= mgh - 0 \\ W &= -mgh \\ \Delta U &= -mgh\end{aligned}$$

$$\text{or } W = \int_I^F \vec{F} \cdot d\vec{s}$$

$$\Delta U = - \int_I^F \vec{F} \cdot d\vec{s}$$

in terms of the electric field

$$\Delta U = -q_0 \int_I^F \vec{E} \cdot d\vec{s}$$

‡ usually, instead of potential energy, we  
 talk about Electric potential,  $V = U/q_0$

‡ as with potential energy, only differences are meaningful

OK, so, we've got this difference in electric potential related to the Electric field

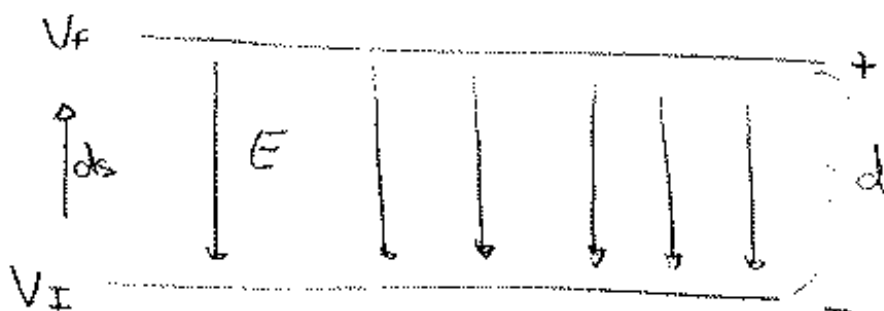
(4)

$$\Delta V = - \int_I^F \vec{E} \cdot d\vec{s}$$

What's so cool about it is that the units for Electric potential are called volts. The same thing you find on a battery. You can use a potential difference to make an electric field! We can't do that very easily with Gravity....

Oscilloscope is a great demonstration of this.

between the plates (Separate them by a distance  $d$ )



higher potential here than here (what would positive charge do?)

$$\Delta V = V_F - V_I = - \int_0^d \vec{E} \cdot d\vec{s} = E_s \Big|_0^d = E_d$$

Confirms what we thought

recall our expression for  $E$ ,  $E = q/\epsilon_0 A$

So, the amount of charge on the plates depends on the  $\Delta V$  & geometry of the conductor;  $q = \frac{\epsilon_0 A}{d} \Delta V$

There are a few tricks that come in handy for electric potential. (5)

1) Electric potential decreases in the direction of the electric field vector

2) Potential differences are meaningful.

3) challenge is to find where to set  $V=0$

ex point charge, what's the potential difference between  $r=\infty$  &  $r=R_0$ ?

$$\begin{aligned}\Delta V = V_F - V_I &= - \int_{R_0}^{\infty} \vec{E} \cdot d\vec{s} = - \int_{R_0}^{\infty} \frac{kq}{r^2} dr \quad \text{for } +q \\ &= - \left( -\frac{kq}{r} \Big|_{R_0}^{\infty} \right) = - \left[ 0 - \left( -\frac{kq}{R_0} \right) \right] \\ &= -\frac{kq}{R_0}\end{aligned}$$

Usually, define  $V_{\infty} = 0$

$$\Delta V = -V_I = -\frac{kq}{R_0} \quad \text{or} \quad V_{at R_0} = \frac{kq}{R_0}$$

&  $V_{at R_0}$  is called the electric potential at  $R_0$  even if it is the difference between  $r=\infty$  &  $R_0$

(easier to do calculations with)

## Example

6



what is the potential at P?



$$V_{\text{tot}} = \sum \frac{kq}{r} = \frac{k(1\mu\text{C})}{.5\text{m}} + \frac{k(1\mu\text{C})}{.5\text{m}} + \frac{k(1\mu\text{C})}{.5\text{m}}$$
$$= \frac{3(8.99 \times 10^9 \text{ N}\frac{\text{m}^2}{\text{C}^2})(10^{-6}\text{C})}{.5\text{m}}$$

$$\sim 5.40 \times 10^4 \frac{\text{Nm}}{\text{C}}$$

lots easier than all those integrals!

4) notice how the integral is set up

$$V = - \int \vec{E} \cdot d\vec{s} = - \int [E_x dx + E_y dy + E_z dz]$$

looks like

$$dV = -E_x dx - E_y dy - E_z dz$$
$$= - \left( \frac{\partial V}{\partial x} \right) dx - \left( \frac{\partial V}{\partial y} \right) dy - \left( \frac{\partial V}{\partial z} \right) dz$$

IF you know V, you can calculate E

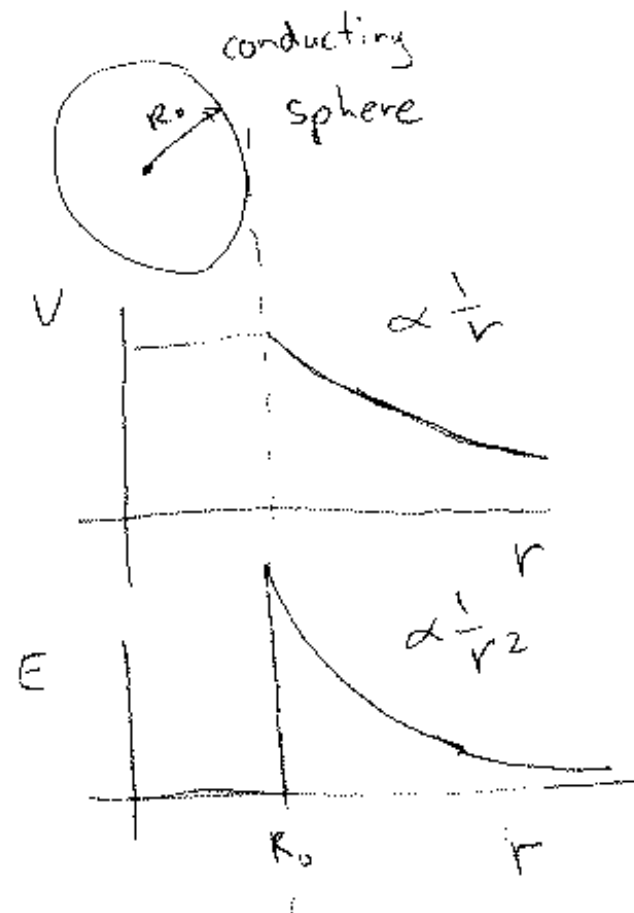
ex We've said  $E=0$  in conductor at equilibrium. Implies that  $V$  is constant

(7)

$$\frac{\partial V}{\partial x} = 0 \text{ inside a conductor}$$

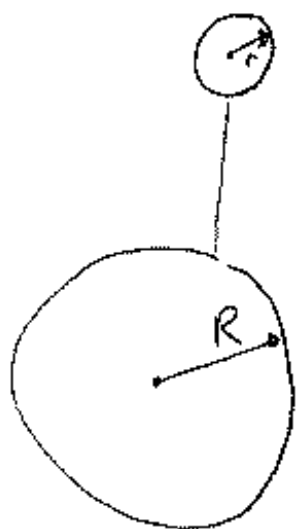
makes sense, recall Franklin found no force on the ball inside the conductor (no change in potential)

Also implies every point on the surface of a conductor is at the same potential as the inside,  $\frac{\partial (\text{constant})}{\partial x} = 0$



This can be very useful!

Suppose we have 2 spherical conductors connected by a thin wire (8)



This means their surfaces are at the same potential.

So, if there's any charge in there

$$\frac{Kq_e}{R} = \frac{Kq_r}{r}$$

What about the electric field

$$\frac{Kq_e}{R^2} = ?$$

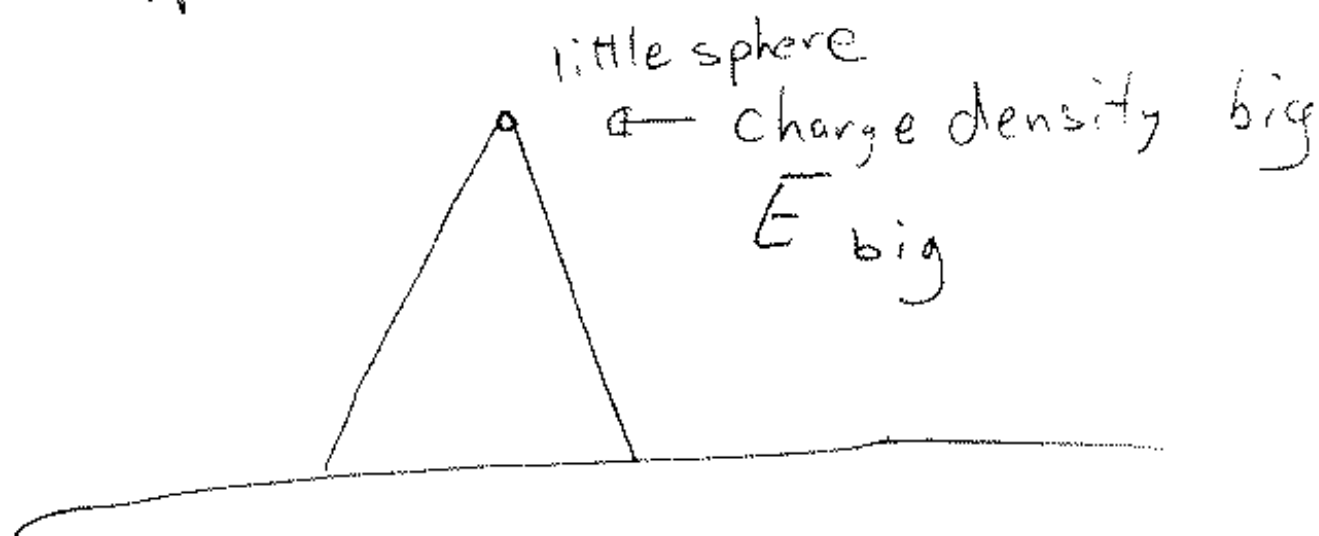
lets look at the ration of electric fields

$$\frac{E_r}{E_R} = \frac{\frac{Kq_r}{r^2}}{\frac{Kq_r}{R^2}} = \frac{\cancel{Kq_r} \frac{1}{r}}{\cancel{Kq_r} \frac{1}{R}} = \frac{\frac{1}{r}}{\frac{1}{R}} = \frac{R}{r}$$

Electric field is bigger on the small sphere = surface charge density bigger on small sphere

Application lightning rods

9



big sphere

Electric field gets so strong  
it starts to break down the  
air (like the little ball & the  
Van De Graff generator)