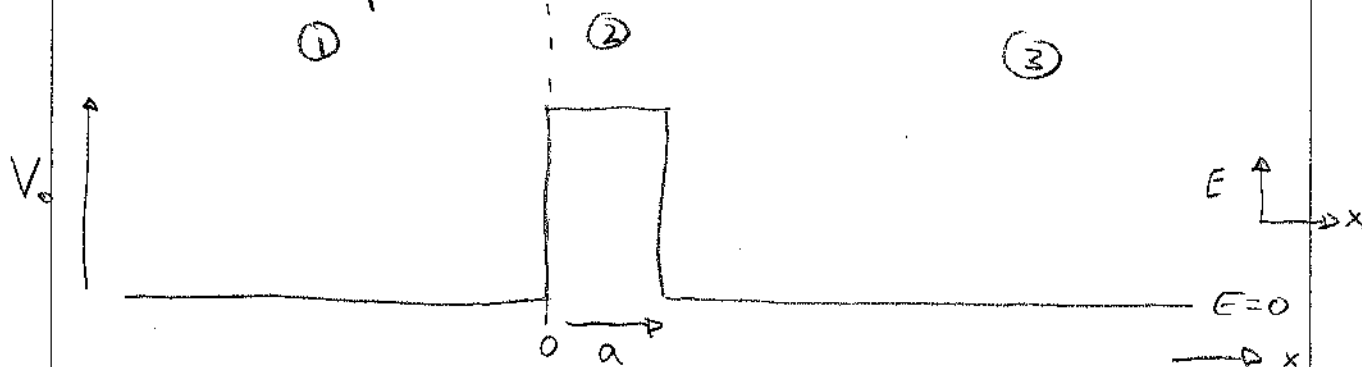


1st, we need to solve the classic "barrier Penetration" problem



In Region ① the wave function has solutions like $Ae^{ik_1x} + Be^{-ik_1x}$

In ② $Ce^{k_2x} + De^{-k_2x}$

③ $Ee^{ik_3x} + Fe^{-ik_3x}$

We can simplify if we only consider waves travelling to the right in region ③ so $F=0$. We can also consider a specific energy of a particle and its de-Broglie wavelength.

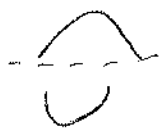
Recall some properties of the wave function

- higher energy - smaller wavelength
- classically forbidden regions - wave functions decay exponentially

$$\frac{d^2\psi}{dx^2} = -\frac{2m}{\hbar^2}(E-V)\psi$$

always towards axis

if $E-V > 0$, $\frac{d^2\psi}{dx^2}$ is negative
positive



if ψ is > 0
if ψ is < 0

always away from axis

if $E-V < 0$ $\frac{d^2\psi}{dx^2}$ is positive
negative



(always)

- Wave functions are smooth

$$\psi_1 = \psi_2, \quad \frac{d\psi_1}{dx} = \frac{d\psi_2}{dx} \quad \text{at } x=0$$

$$\psi_2 = \psi_3, \quad \frac{d\psi_2}{dx} = \frac{d\psi_3}{dx} \quad \text{at } x=a$$

$$A+B = C+D$$

$$(Aik_1 - Bik_1) = k_2(C-D)$$

$$Ee^{ik_3a} = Ce^{k_2a} + De^{-k_2a}$$

$$ik_3 Ee^{ik_3a} = k_2 (Ce^{k_2a} - De^{-k_2a})$$

$$k_1 = \sqrt{2mE}/\hbar$$

$$k_2 = \sqrt{2m(U-E)}/\hbar$$

$$k_3 = \sqrt{2mE}/\hbar = k_1$$

yikes! These are easier to solve if we let $E > U$ (st \ddagger) then relax this condition

we get:

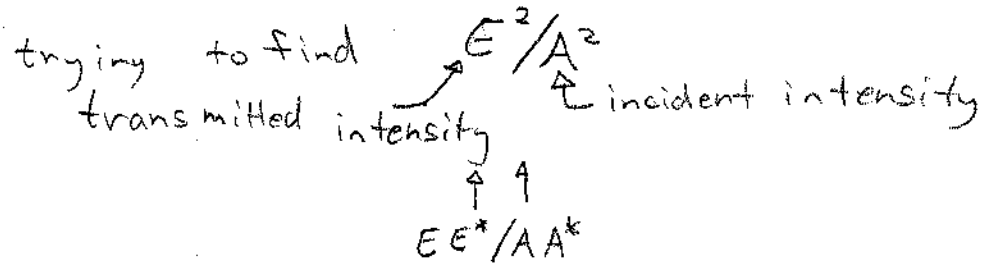
$$Aik_1 - Bik_2 = ik_2(C-D) = ik_1(A-B)$$

$$A+B = C+D$$

$$Ee^{ik_1a} = Ce^{ik_2a} + De^{-ik_2a}$$

$$Ek_1 e^{ika} = k_2 (Ce^{ik_2a} - De^{-ik_2a})$$

5 equations \ddagger 4 unknowns - traditionally we leave A as a variable and look at quantities in relation to A. In this case we are



$$k_1(A+B) + k_1(A-B) = k_1(C+D) + k_2(C-D)$$

$$2k_1A = (k_1+k_2)C + (k_1-k_2)D$$

$$k_2 E e^{ik_1 a} + k_1 E e^{ik_1 a} = k_2 C e^{ik_2 a} + k_2 D e^{ik_2 a} + k_2 C e^{ik_2 a} - k_2 D e^{-ik_2 a}$$

$$(k_2+k_1) E e^{ik_1 a} = 2k_2 C e^{ik_2 a}$$

$$(k_2-k_1) E e^{ik_1 a} = 2k_2 D e^{-ik_2 a}$$

$$\text{or } C = \frac{(k_2+k_1)}{2k_2} E e^{ik_1 a} e^{-ik_2 a}$$

$$D = \frac{(k_2-k_1)}{2k_2} E e^{ik_1 a} e^{ik_2 a}$$

$$\xi \quad 2k_1 A = \frac{(k_1+k_2)^2}{2k_2} E e^{ik_1 a} e^{-ik_2 a} - \frac{(k_2-k_1)^2}{2k_2} E e^{ik_1 a} e^{ik_2 a}$$

$$A = \frac{(k_1+k_2)^2}{4k_1 k_2} E e^{ik_1 a} e^{-ik_2 a} - \frac{(k_2-k_1)^2}{4k_1 k_2} E e^{ik_1 a} e^{ik_2 a}$$

substitute $k_2 = i\alpha$ (for $E < V$)

$$\frac{AA^*}{EE^*} = \left(\frac{(k_1+i\alpha)^2}{4k_1 i\alpha} e^{ik_1 a} e^{\alpha a} - \frac{(k_1-i\alpha)^2}{4k_1 i\alpha} e^{ik_1 a} e^{-\alpha a} \right) \times$$

$$\left(\frac{(k_1-i\alpha)^2}{-4k_1 i\alpha} e^{-ik_1 a} e^{\alpha a} + \frac{(k_1+i\alpha)^2}{4k_1 i\alpha} e^{-ik_1 a} e^{-\alpha a} \right)$$

$$= \frac{(k_1^2 + \alpha^2)^4}{(4k_1 \alpha)^2} e^{2\alpha a} + \frac{(k_1^2 + \alpha^2)^4}{(4k_1 \alpha)^2} e^{-2\alpha a} - \frac{(k_1+i\alpha)^4}{(4k_1 \alpha)^2} - \frac{(k_1-i\alpha)^4}{(4k_1 \alpha)^2}$$

note: $(k_1+i\alpha)^4 = k_1^4 - 4i\alpha k_1^3 - 6k_1^2 \alpha^2 + 4i\alpha^3 k_1 + \alpha^4$

$(k_1-i\alpha)^4 = k_1^4 + 4i\alpha k_1^3 - 6k_1^2 \alpha^2 - 4i\alpha^3 k_1 + \alpha^4$

$$(k_1+i\alpha)^4 + (k_1-i\alpha)^4 = 2(k_1^4 - 6k_1^2 \alpha^2 + \alpha^4) = 2(k_1^2 + \alpha^2)^2 - 16k_1^2 \alpha^2$$

$$\frac{AA^*}{EE^*} = \left[\frac{(k_1^2 + \alpha^2)(e^{\alpha a} - e^{-\alpha a})}{2k_1 \alpha} \right]^2 + \frac{16k_1^2 \alpha^2}{16k_1^2 \alpha^2}$$

$$T = \frac{AA^*}{EE^*} = \left[\frac{(2mE + (2m)(V-E))^2}{4(2mE)(2m)(V-E)} \sinh^2(\alpha a) + 1 \right]$$

$$= \left[\frac{V^2}{4(V-E)E} \sinh^2(\alpha a) + 1 \right]$$

Transmission probability = $\left[1 + \frac{\sinh^2(\alpha a)}{4 \frac{E}{V} (1 - \frac{E}{V})} \right]^{-1}$

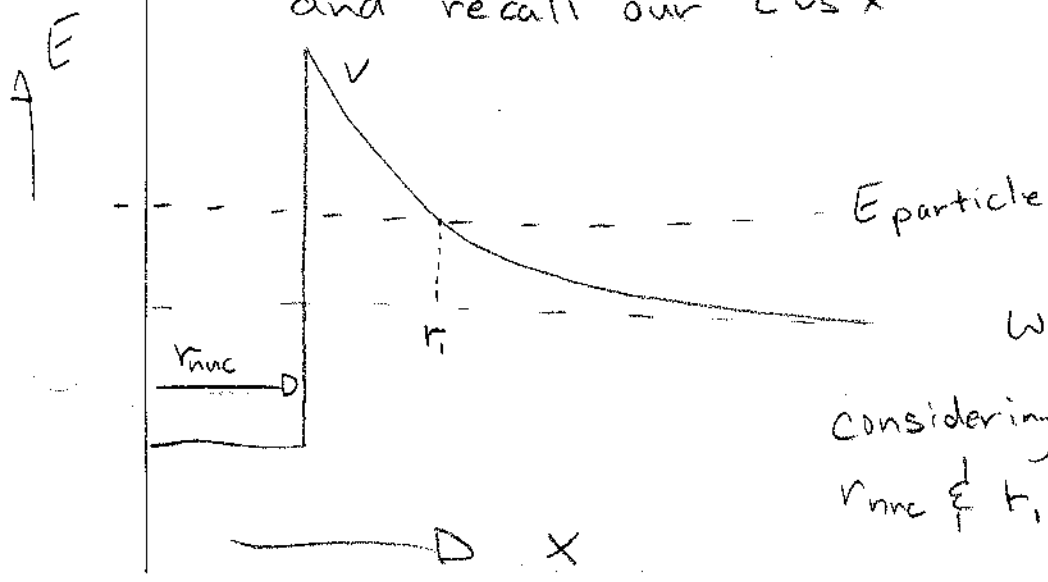
For our alpha decay, we will treat the transmission probability as small with $\sinh^2(\alpha a)$ large

so $T = 16 \frac{E}{V} (1 - \frac{E}{V}) e^{-2\alpha a}$

To get a feel for the order of magnitude behavior of this, notice, for "reasonable" values of $E \ll V$ ($E \sim 4 \text{ MeV}$, $V \sim 30 \text{ MeV}$) $16 \frac{E}{V} (1 - \frac{E}{V})$ is essentially $\propto 1$ over the range $\frac{1}{10} = \frac{E}{V}$ to nearly $1 = \frac{E}{V}$, so we'll neglect it. Now consider the remaining term

$$T \sim e^{-2\alpha a}$$

and recall our E vs x



We proceed by considering T between $r_{nuc} \& x_1$